## Chapter 18

## A First Look at Stochastic Integrals with the Wiener Process

Section 18.1 touches briefly on the martingale characterization of the Wiener process.

Section 18.2 gives a heuristic introduction to stochastic integrals, via Euler's method for approximating ordinary integrals.

## 18.1 Martingale Characterization of the Wiener Process

Because the Wiener process is a Lévy process (Example 139), it is self-similar in the sense of Definition 147. That is, for any  $a > 0$ ,  $W(at) \stackrel{d}{=} a^{1/2}W(t)$ . In fact, if we define a new process  $W_a$  through  $W_a(t, \omega) = a^{-1/2}W(at, \omega)$ , then  $W_a$  is itself a Wiener process. Thus the whole process is self-similar. This is only one of several sorts of self-similarities in the Wiener process. Another is sometimes called *spatial homogeneity*:  $W_{\tau}$ , defined through  $W_{\tau}(t,\omega) = W(t +$  $\tau, \omega$ ) –  $W(\tau, \text{omega})$  is also a Wiener process. That is, if we "re-zero" to the state of the Wiener process  $W(\tau)$  at an arbitrary time  $\tau$ , the new process looks just like the old process.  $-W(t)$ , obviously, is also a Wiener process.

Related to these properties is the fact that  $W^2(t) - t$  is a martingale with respect to  $\{\mathcal{F}_t^W\}$ . (This is easily shown with a little algebra.) What is more surprising is that this is enough to characterize the Wiener process.

Theorem 224 (Martingale Characterization of the Wiener Process) *If*  $M(t)$  *is a continuous martingale, and*  $M^2(t) - t$  *is also a martingale, then*  $M(t)$ *is a Wiener process.*

There are some very clean proofs of this theorem<sup>1</sup> — but they require us to use stochastic calculus! Doob (1953, pp. 384ff) gives a proof which does not, however. The details of his proof are messy, but the basic idea is to get the central limit theorem to apply, using the martingale property of  $M^2(t) - t$  to get the variance to grow linearly with time and to get independent increments, and then seeing that between any two times  $t_1$  and  $t_2$ , we can fit arbitrarily many little increments so we can use the CLT.

We will return to this result as an illustration of the stochastic calculus (Theorem 247).

## 18.2 A Heuristic Introduction to Stochastic Integrals

Euler's method is perhaps the most basic method for numerically approximating integrals. If we want to evaluate  $I(x) \equiv \int_a^b x(t)dt$ , then we pick n intervals of time, with boundaries  $a = t_0 < t_1 < \ldots t_n = b$ , and set

$$
I_n(x) = \sum_{i=1}^n x(t_{i-1})(t_i - t_{i-1})
$$

Then  $I_n(x) \to I(x)$ , if x is well-behaved and the length of the largest interval  $\to 0$ . If we want to evaluate  $\int_{t=a}^{t=b} x(t)dw$ , where w is another function of t, the natural thing to do is to get the derivative of  $w, w'$ , replace the integrand by  $x(t)w'(t)$ , and perform the integral with respect to t. The approximating sums are then

$$
\sum_{i=1}^{n} x(t_{i-1}) w'(t_{i-1}) (t_i - t_{i-1})
$$
\n(18.1)

Alternately, we could, if  $w(t)$  is nice enough, approximate the integral by

$$
\sum_{i=1}^{n} x(t_{i-1}) (w(t_i) - w(t_{i-1}))
$$
\n(18.2)

even if  $w'$  doesn't exist.

(You may be more familiar with using Euler's method to solve ODEs,  $dx/dt =$  $f(x)$ . Then one generally picks a  $\Delta t$ , and iterates

$$
x(t + \Delta t) = x(t) + f(x)\Delta t \tag{18.3}
$$

from the initial condition  $x(t_0) = x_0$ , and uses linear interpolation to get a continuous, almost-everywhere-differentiable curve. Remarkably enough, this converges on the actual solution as  $\Delta t$  shrinks (Arnol'd, 1973).)

Let's try to carry all this over to random functions of time  $X(t)$  and  $W(t)$ . The integral  $\int X(t)dt$  is generally not a problem — we just find a version of X

<sup>&</sup>lt;sup>1</sup>See especially Ethier and Kurtz (1986, Theorem 5.2.12, p. 290).

with measurable sample paths (Section 8.2).  $\int X(t)dW$  is also comprehensible if  $dW/dt$  exists (almost surely). Unfortunately, we've seen that this is not the case for the Wiener process, which (as you can tell from the  $W$ ) is what we'd really like to use here. So we can't approximate the integral with a sum like Eq. 18.1. But there's nothing preventing us from using one like Eq. 18.2, since that only demands increments of  $W$ . So what we would like to say is that

$$
\int_{t=a}^{t=b} X(t)dW \equiv \lim_{n \to \infty} \sum_{i=1}^{n} X(t_{i-1}) (W(t_i) - W(t_{i-1})) \tag{18.4}
$$

This is a crude-but-workable approach to numerically evaluating stochastic integrals, and apparently how the first stochastic integrals were defined, back in the 1920s. Notice that it is going to make the integral a *random variable*, i.e., a measurable function of  $\omega$ . Notice also that I haven't said anything yet which should lead you to believe that the limit on the right-hand side exists, in any sense, or that it is independent of the choice of partitions  $a = t_0 < t_1 < \dots t_n$  b. The next chapter will attempt to rectify this.

(When it comes to the SDE  $dX = f(X)dt + g(X)dW$ , the counterpart of Eq. 18.3 is

$$
X(t + \Delta t) = X(t) + f(X(t))\Delta t + g(X(t))\Delta W \tag{18.5}
$$

where  $\Delta W = W(t+\Delta t) - W(t)$ , and again we use linear interpolation in between the points, starting from  $X(t_0) = x_0$ .